

CHEMICAL REACTION EFFECTS ON MHD FREE CONVECTION HEAT AND MASS TRANSFER FLOW THROUGH A POROUS MEDIUM BOUNDED BY TWO VERTICAL WALLS IN THE PRESENCE OF RADIATION

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Abstract: *The present paper analyzes the effect of chemical reaction on free convective magneto hydrodynamic (MHD) flow of steady, laminar, incompressible fluid between two vertical walls filled with porous materials. The flow passes through an exponential radiative stretching sheet in the presence of magnetic field. The governing PDE of the fluid flow are solved using perturbation technique. Approximate solutions are derived for velocity, temperature and concentration field. The influences of various physical parameters on velocity, temperature, and solutal concentration profiles are presented through graphs and the numerical computation of physical quantities such as skin friction coefficient, Sherwood number, Nusselt number are presented and discussed.*

1.1 Introduction

Problems involving Magneto Hydrodynamics (MHD) are very important in many fields such as geophysical and astrophysical problems, plasma studies, nuclear reactors, geothermal energy extractions and the boundary layer control in the field of aerodynamics. The possible usage of MHD is effect a flow stream of an electrically conducting fluid for the purpose of thermal production, braking, propulsion and control. The effects of Magnetic field on free convection flow problems have been attracted by many investigators [6-8]. Chemical reactions can be codified as either heterogeneous or homogeneous processes. This depends on whether they occur at an interface or as a single phase volume reaction. In well-mixed systems, the reaction is heterogeneous, if it takes place at an interface and homogeneous, if it takes place in solution. In most cases of chemical reactions, the reaction rate depends on the concentration of the species itself.

Chemical reaction have plentiful industrial applications such as several engineering, industrial, astrophysical and geophysical application, such as polymer production, manufacturing of ceramic, packed-bed catalytic reactors, enhanced oil recovery, food processing, underground energy transport, cooling of nuclear reactors, high-speed plasma wind, magnetized plasma flow, cosmic jets and stellar systems. A clear understanding of the nature of interaction between thermal and concentration buoyancies is necessary to control these processes. Anjali Devi and Kandasamy 2000 investigated the effects of chemical reaction heat and mass transfer on MHD flow past a semi-infinite plate. Chamkha. J Ali, (2004) studied "heat and mass transfer past a semi-infinite vertical permeable moving plate with heat absorption Kandaswamy.R , Periasamy K, SivagnanaPrabhu KK (2005), studied Chemical reaction, heat and mass transfer on MHD flow over a vertical stretching surface with heat source and thermal stratification effects. Murti et al 2005, studied the Radiation, chemical reaction, Double dispersion effects on Heat and mass transfer in non-Newtonian fluids. Muthukumaraswamy 2010 studied Chemical Reaction Effects on Vertical Oscillating Plate with Variable Temperature.

Radiation heat transfer is ubiquitous, because all matter emits and absorbs electromagnetic radiation. The electromagnetic radiation spectrum is huge, but heat transfer is mostly concerned with a small part of

it, called thermal radiation. Thermal radiation effect plays a significant role in controlling heat transfer procession polymer processing industry. The quality of the final product depends to a certain extent on heat controlling factors. Also, the effects of thermal radiation on flow and heat transfer processes are important in the design of many advanced energy convection systems which operate at high temperature. Thermal radiation occurring within these systems is usually the result of emission by the hot walls and the working fluid. Thermal radiation effects become more important when the difference between the surface and the ambient temperature is large. Thus thermal radiation is one of the vital factors controlling the heat and mass transfer. The knowledge of radiation heat transfer in the system can perhaps lead to a desired product with sought characteristics. Pal 2009, analyzed heat and mass transfer in two - dimensional stagnation-point flow of an incompressible viscous fluid over a stretching vertical sheet in the presence of buoyancy force and thermal radiation. Pal and Mondal 2009, studied the influence of temperature - dependent viscosity and thermal radiation on MHD forced convection over a non-isothermal wedge.

1.2 Mathematical formulation

We have considered a two dimensional steady flow of a laminar free convective incompressible electrically conducting fluid past an infinite vertical porous plate, embedded in a porous medium in the presence of thermal radiation and a homogeneous chemical reaction. A uniform magnetic field of strength B_0 is applied perpendicular to the plate. Let h be the distance between the plates. It is assumed that there is applied voltage which results the absence of electric field. The transversely applied magnetic field and magnetic Reynolds number are assumed to be very small so that the induced magnetic field and Hall currents are negligible [6]. The suction velocity is assumed in the form $v = -v_0$, where v_0 is the constant suction /injection.

$$\frac{\partial v}{\partial y} = 0 \tag{1.1}$$

$$v \frac{\partial u}{\partial y} = v \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty) + g\beta^*(C - C_\infty) - \left(\frac{\sigma B_0^2}{\rho} \right) u + \left(\frac{v}{k} \right) u \tag{1.2}$$

$$v \frac{\partial T}{\partial y} = \frac{\lambda}{\rho C_p} \frac{\partial^2 T}{\partial y^2} - \frac{1}{\rho C_p} \frac{\partial q_r}{\partial y} - \frac{Q_0}{\rho C_p} (T - T_\infty) + Q_1'(C - C_\infty) \tag{1.3}$$

$$v \frac{\partial C}{\partial y} = D \frac{\partial^2 C}{\partial y^2} - Kr'(C - C_\infty) \tag{1.4}$$

Where u and v are the velocity components along and perpendicular to the surface, g is the acceleration due to gravity, T the temperature of the fluid near the plate, T_∞ the free stream temperature, C concentration, β the coefficient of thermal expansions, β^* the volumetric coefficient of expansion of the species concentration, k the thermal conductivity, C_p the specific heat of constant pressure, B_0 the magnetic field coefficient, μ viscosity of the fluid, the density, σ the magnetic permeability of fluid, V_0 constant suction velocity, ν the kinematic viscosity and D chemical molecular diffusivity.

The relevant boundary conditions are

$$u = 0, \quad \theta = 1, C = 1 \quad \text{at } y = 0$$

$$u = 0, \quad \theta = 0, C = 0 \quad \text{at } y = 1 \quad (1.5)$$

$$v = -v_0 \text{ constant.}$$

Where $v_0 > 0$ corresponds to steady suction velocity normal.

The local radiant for the case of an optically thin gray gas is expressed by

$$\frac{\partial q_r}{\partial y} = -4\alpha^* \sigma (T_\infty^4 - T^4) \quad (1.6)$$

It is assume that the temperature differences within the flow are sufficiently small such that T^4 may be expressed as a linear function of the temperature. This is accomplished by expanding T^4 in a Taylor series about T_∞ and neglecting higher-order terms, thus

$$T^4 \cong 4T_\infty^3 T - 3T_\infty^4 \quad (1.7)$$

By using equations (1.6) and (1.7), equation (1.3) reduces to

$$-v_0 \frac{\partial T}{\partial y} = \frac{\lambda}{\rho C_p} \frac{\partial^2 T}{\partial y^2} + 16\alpha^* \sigma T_\infty^3 (T - T_\infty) - \frac{Q_0}{\rho C_p} (T - T_\infty) + Q_1' (C - C_\infty) \quad (1.8)$$

We introduce the following non-dimensional quantities

$$u = \frac{u^*}{v_0}, \quad y = \frac{v_0 y^*}{v}, \quad M = \frac{\sigma B_0^2 v}{v_0^2 \rho}, \quad \theta = \frac{(T - T_\infty) v_0 \lambda}{qv}, \quad \text{Pr} = \frac{\mu C_p}{\lambda}$$

$$\alpha = \frac{V_0^2 K}{v^2}, \quad C = \frac{(C - C_\infty) v_0 D}{mv}, \quad Gr = \frac{\rho \beta g v^2 (T_w - T_\infty)}{v_0^3 \mu}, \quad Gc = \frac{\rho \beta g (C - C_\infty)}{v_0^3}$$

$$Sc = \frac{v}{D}, \quad Kr = \frac{Kr' m v}{D V_0^2}, \quad Q_0 = \frac{\phi V_0^2 \rho C_p}{v}, \quad Q_1' = \frac{Q_1 V_0^2 D}{m v \lambda}, \quad R = \frac{16 \alpha v^2 \sigma T_\infty^3}{k v_0^2} \quad (1.9)$$

Where Gr is Grashof number, Pr is Prandtl number, M is Magnetic number, Sc is Schmidt number, Kr is Chemical reaction parameter, Q_1 is heat absorption parameter, ϕ is heat source parameter, and R Radiation parameter where q is the heat flux term per unit area and m is the mass flux per unit area.

We get

$$-\frac{\partial u}{\partial y} = \frac{\partial^2 u}{\partial y^2} + Gr\theta + GcC - (\alpha^{-1} + M)u \quad (1.10)$$

$$-\frac{\partial \theta}{\partial y} = \frac{1}{Pr} \frac{\partial^2 \theta}{\partial y^2} - R\theta - \theta\phi - Q_1 C \quad (1.11)$$

$$-\frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} - KrC \quad (1.12)$$

Subject to the boundary conditions

$$u = 0, \quad \theta = 1, \quad C = 1 \quad \text{at} \quad y = 0$$

$$u = 0, \quad \theta = 0, \quad C = 0 \quad \text{at} \quad y = 1 \quad (1.13)$$

1.3 Solution of the problem

In order to solve the set partial differential equations (1.10) to (1.12) in non-dimensional form subject to boundary conditions, we assumed the velocity, temperature and concentration in a series expansion in powers of E where $E \ll 1$ as given below.

$$u(y) = u_0(y) + Eu_1(y) + O(\varepsilon^2) \quad (1.14)$$

$$\theta(y) = \theta_0(y) + E\theta_1(y) + O(\varepsilon^2)$$

$$C(y) = C_0(y) + EC_1(y) + O(\varepsilon^2)$$

Substituting equations (1.14) into equations (1.10) to (1.12) and equating the coefficient of similar powers of E and neglecting the higher powers of E , we obtain the following ordinary differential equations for (u_0, θ_0, C_0) and

(u_1, θ_1, C_1) .

$$u_0'' + u_0' - (\alpha^{-1} + M)u_0 = -Gr\theta_0 - GcC_0 \quad (1.15)$$

$$\theta_0'' + Pr\theta_0' - Pr(R + \phi)\theta_0 = Pr Q_1 C_0 \quad (1.16)$$

$$C_0'' + ScC_0' - KrScC_0 = 0 \quad (1.17)$$

$$u_1'' + u_1' - (\alpha^{-1} + M)u_1 = -Gr\theta_1 - GcC_1 \quad (1.18)$$

$$\theta_1'' + Pr \theta_1' - Pr(R + \phi)\theta_1 = Pr Q_l C_1 \quad (1.19)$$

$$C_1'' + Sc C_1' - Kr Sc C_1 = 0 \quad (1.20)$$

and the corresponding boundary conditions are

$$\begin{aligned} u_0 = 0, \quad u_1 = 0, \quad \theta_0 = 1, \quad \theta_1 = 0, \quad C_0 = 1, \quad C_1 = 0 \quad \text{at } y = 0 \\ u_0 = 0, \quad u_1 = 0, \quad \theta_0 = 0, \quad \theta_1 = 0, \quad C_0 = 0, \quad C_1 = 0 \quad \text{at } y = 1 \end{aligned} \quad (1.21)$$

Solving equations (1.15) to (20).with the help of (1.21)., we get

$$u_0 = P_6 e^{m_5 y} + P_5 e^{-m_6 y} + P_1 e^{m_3 y} + P_2 e^{-m_4 y} + P_3 e^{m_1 y} + P_4 e^{-m_2 y}$$

$$u_1 = 0$$

$$\theta_0 = B_4 e^{m_3 y} + B_3 e^{-m_4 y} + A_1 e^{m_1 y} + A_2 e^{-m_2 y}$$

$$\theta_1 = 0$$

$$C_0 = B_1 e^{m_1 y} + B_2 e^{-m_2 y}$$

$$C_1 = 0$$

$$u(y) = P_6 e^{m_5 y} + P_5 e^{-m_6 y} + P_1 e^{m_3 y} + P_2 e^{-m_4 y} + P_3 e^{m_1 y} + P_4 e^{-m_2 y}$$

$$\theta(y) = B_4 e^{m_3 y} + B_3 e^{-m_4 y} + A_1 e^{m_1 y} + A_2 e^{-m_2 y}$$

$$C(y) = B_1 e^{m_1 y} + B_2 e^{-m_2 y}$$

The skin-friction coefficient, the Nusselt number and the Sherwood number near the plate, are important physical parameters for this type of boundary-layer flow. These parameters can be defined and determined as follows:

The skin-friction coefficient at the plate is given by

$$\tau = \left(\frac{\partial u}{\partial y} \right)_{y=0} = m_5 P_6 - m_6 P_5 + m_3 P_1 - m_4 P_2 + m_1 P_3 - m_2 P_4$$

The rate of heat transfer in terms of Nusselt number at the plate is given by

$$Nu = \left(\frac{\partial \theta}{\partial y} \right)_{y=0} = m_3 B_4 - m_4 B_3 + m_1 A_1 - m_2 A_2$$

The Sherwood number at the plate is given by

$$Sh = \left(\frac{\partial C}{\partial y} \right)_{y=0} = m_1 B_1 - m_2 B_2$$

1.4 Results and Discussions

The numerical values of the velocity, temperature, concentration, Skin friction, Nusselt number and Sherwood number are computed for different values of physical parameters like magnetic parameter (M), radiation parameter (R), chemical reaction (Kr) the effects of these parameters on flow quantities are studied through graphs 1.1-1.16.

Figure 1.1 plots the velocity profiles against the span-wise coordinate y for different magnetic field parameter M . This illustrates that velocity decreases as the existence of magnetic field becomes stronger. This conclusion agrees with the fact that magnetic field exerts retarding force on the free-convection flow.

From figure 1.2 it is seen, that the velocity decreases with increasing the Chemical reaction parameter Kr .

Figure 1.3 illustrates the dimensionless velocity u for different values of the Prandtl number Pr . The analytical results show that the effect of increasing values of Prandtl number results in a decreasing velocity. It is observed from Figure 1.4 that an increase in Grashof number for heat transfer Gr leads to a rise in the values of velocity u due to enhancement in buoyancy force. The plot of velocity profile for different values of Grashof number for mass transfer Gc is given in Figure 1.5. It is observed that velocity increases for the increasing values of Grashof number for mass transfer Gc .

The effect of the Radiation parameter R on the dimensionless velocity u is shown in Figure 6. Figure 1.6 shows that velocity component decreases with an increase in the radiation parameter R . The influence of the Schmidt number Sc on velocity profiles are plotted in Figure 1.7. The Schmidt number Sc embodies the ratio of the momentum to the mass diffusivity. It is noticed that as the Schmidt number Sc increases the velocity decreases. This causes the concentration buoyancy effects to decrease yielding a reduction in the fluid velocity.

Figure 1.8 depicts the dimensionless velocity component u profiles for different values of heat sink parameter Q_l . It is noticed that an increase in the heat sink parameter Q_l results in a decrease in the dimensionless velocity component u within the boundary layer.

The concentration profiles for different values of Schmidt number Sc are plotted in Figure 1.11. For Schmidt number Sc , the value 0.6 corresponds to water vapor and represents a diffusing chemical species of most common interest in air. The analytical results show that the effect of increasing Schmidt number in air and water results in a decreasing concentration distribution across the boundary layer. The results also indicate that the effect on Schmidt number on the velocity and concentration is more in water as compared to that in air.

Figure 1.12 displays that the concentration decreases with increasing the Chemical reaction parameter Kr .

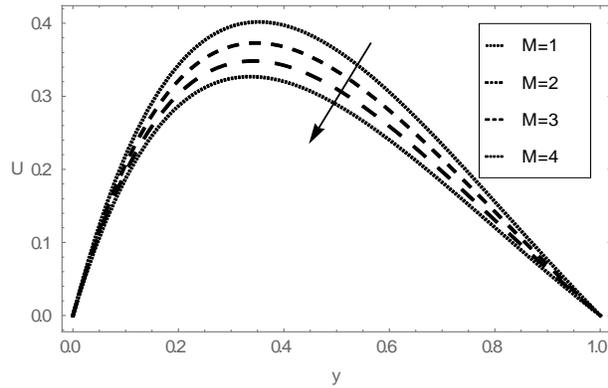


Fig.1.1 Velocity profiles for various values of M

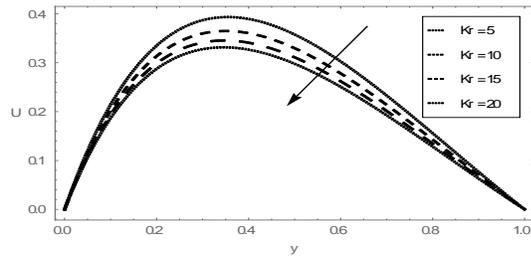


Fig.1.2 Velocity profiles for various values of Kr

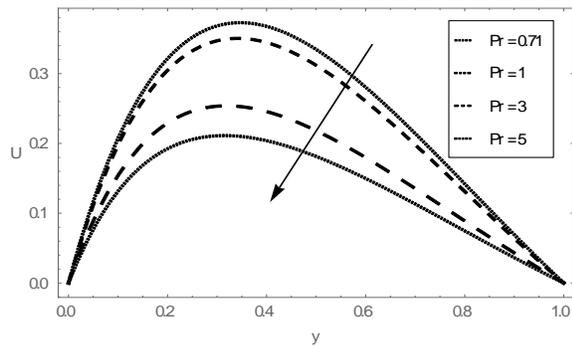


Fig.1.3 Velocity profiles for various values of Pr

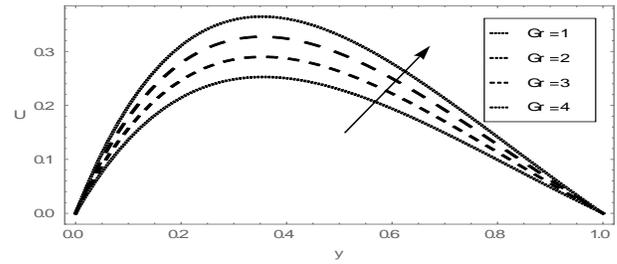


Fig.1.4 Velocity profiles for various values of Gr

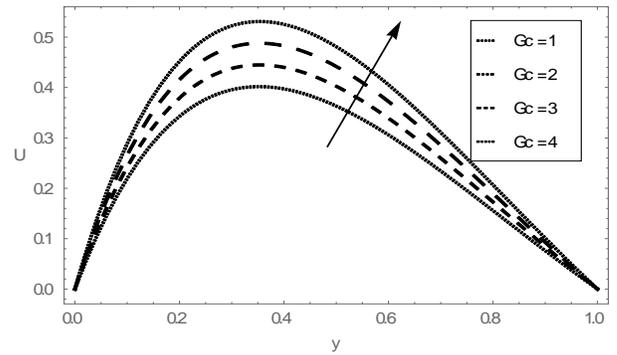


Fig.1.5 Velocity profiles for various values of Gc

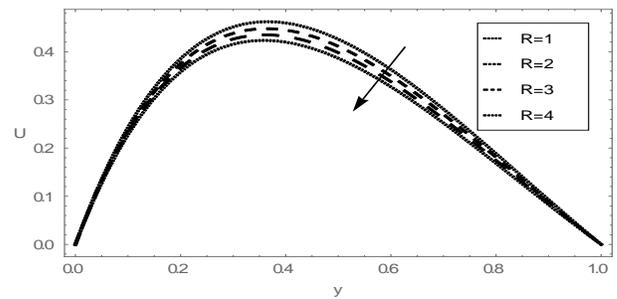


Fig.1.6 Velocity profiles for various values of R

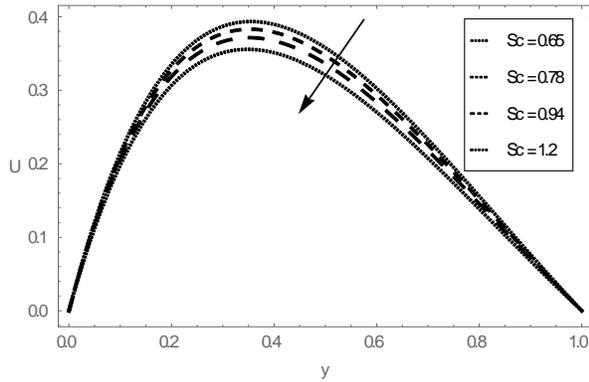


Fig.1.7 Velocity profiles for various values of Sc

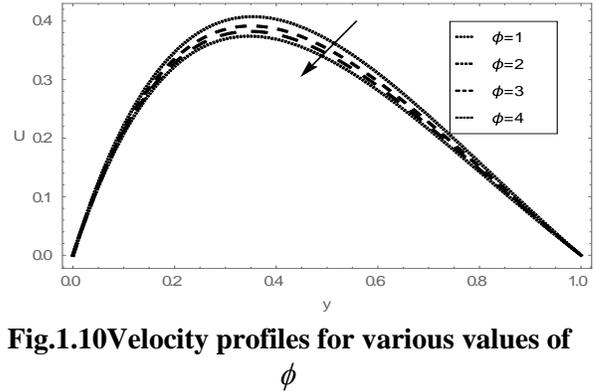


Fig.1.10 Velocity profiles for various values of ϕ

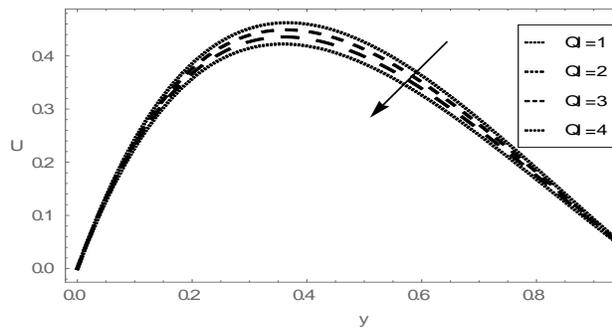


Fig.1.8 Velocity profiles for various values of Q_1

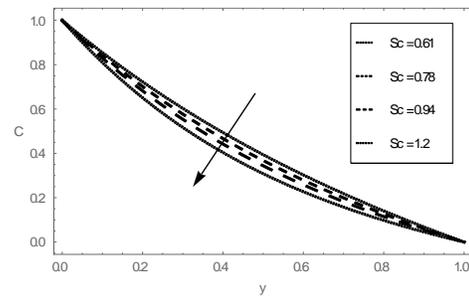


Fig.1.11 concentration profiles for various values of Sc

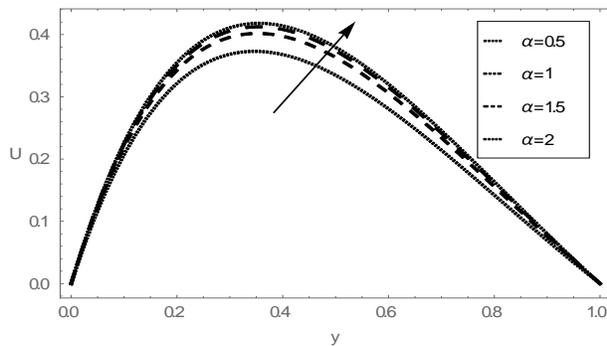


Fig.1.9 Velocity profiles for various values of α

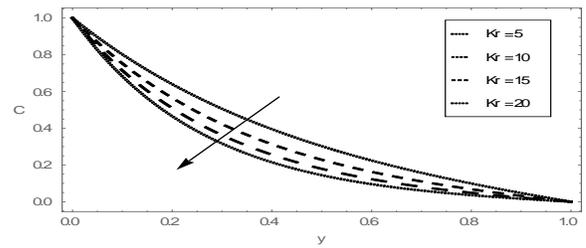


Fig.1.12 concentration profiles for various values of Kr

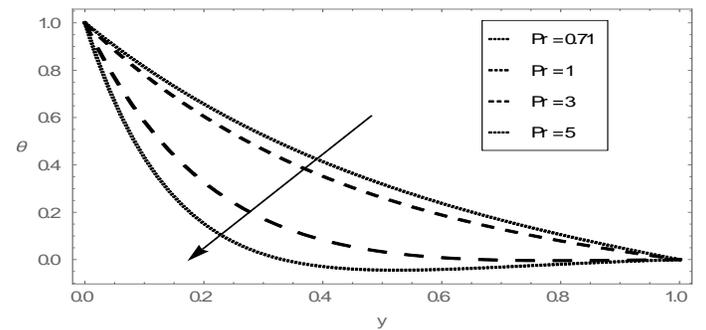


Fig 1.13 Temperature profiles for various values of Pr

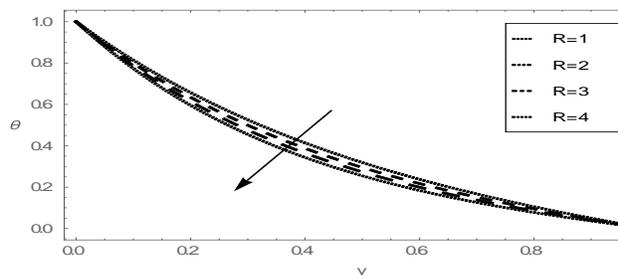


Fig.1.14 Temperature profiles for various values of R

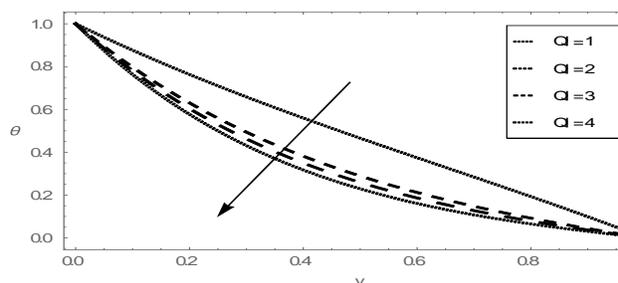


Fig.1.15 Temperature profiles for various values of Q1

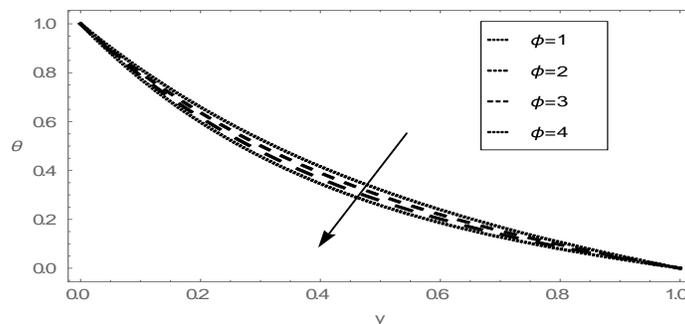


Fig1.16 Temperature profiles for various values of phi

From figure 1.13, it is observed that an increase in the Prandtl number results in a decrease of the thermal boundary layer thickness and in general a lower average temperature within the boundary layer. The reason is that smaller values of Pr are equivalent to increasing the thermal conductivities, and therefore heat is able to diffuse away from the heated plate more rapidly than for higher values of Pr . Hence in the case of smaller Prandtl numbers as the boundary layer is thicker and the rate of heat transfer is reduced.

Figure 1.14 shows the variation of temperature profiles with respect to the radiation parameter R . From this figure, it is observed that as temperature decreases for the increasing values of radiation parameter R . This result qualitatively agrees with expectations, since the effect of radiation is to decrease the rate of energy transport to the fluid, thereby decreasing the temperature of the fluid. The influence of the parameter Q_1 on dimensionless temperature profiles θ is plotted in Figure 1.15. It is noticed that dimensionless temperature decreases with an increase in heat sink parameter Q_1 .

Tables 1.1-1.5 show the effects of the thermal Grashof number, radiation parameter, chemical reaction parameter, the Schmidt number and magnetic parameter on skin friction C_f , the Nusselt number Nu , and Sherwood number Sh . From table 1.1 it is observed that as Gr increased skin friction increased. From table 1.2 it can be easily seen that as radiation parameter increases radiation parameter increases, the skin friction decreases and Nusselt number decreases. From table 1.3 it is noticed that as Kr increases the skin friction decreases and the Sherwood number decreases. From table 1.4 it is observed that as Schmidt number increases the skin friction and Sherwood number decreases. Finally table 1.5 shows that as M increases the skin friction decreases.

Table 1.1 Effects of Gr on skin friction

Gr	Cf
1	1.71797
2	1.9828

1	1.71797
2	1.9828

3	2.24764
4	2.51247

Table 1.2 Effects of radiation on skin friction and Nusselt number

R	Cf	Nu
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Table 1.3 Effects of Kr on skin friction and Sherwood number

Kr	Cf	Sh
5	2.60269	-2.18318
10	2.45868	-2.82813
15	2.35769	-3.35917
20	2.28196	-3.81746S

Table 1.4 Effects of Sc on skin friction and Sherwood number

Sc	Cf	Sh
0.66	2.77465	-2.27064

1	2.77465	-2.03112
2	2.72817	-2.21323
3	2.6864	-2.38345
4	2.64861	-2.54346

0.78	2.74324	-2.47614
0.94	2.70254	-2.74145
1.2	2.63933	-3.115464

Table 1.5 Effects of M on skin friction

M	Cf
1	2.77731
2	2.64626
3	2.52998
4	2.42986

1.5. Conclusions

The governing equations for MHD free convection flow through porous medium bounded by two vertical porous plates in presence of radiation and chemical reaction is formulated. The resulting partial differential equations were transformed into a set of ordinary differential equations using a two term series and solved in closed form. Numerical computations of the closed form results are performed and some graphical results were obtained to illustrate the details of the flow and heat and mass transfer characteristics and their dependence on some of the physical parameters. It is found that velocity decreases with the increase in M, Kr and Pr. Temperature decreases with the increasing values of Pr. However concentration decreases with an increase in Kr.

1.6 References

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